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# Discovery of Plutonium-Based Superconductivity

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**Abstract.** The discovery of superconductivity in single crystals of  $\text{PuCoGa}_5$  with transition temperature  $T_c=18.5$  K is discussed. The existing data lead to the speculation that the superconductivity in  $\text{PuCoGa}_5$  may be unconventional. In such a scenario the properties of  $\text{PuCoGa}_5$  would be intermediate between those of isostructural  $\text{UCoGa}_5$  and  $\text{CeCoIn}_5$ , more heavily studied f-electron materials.

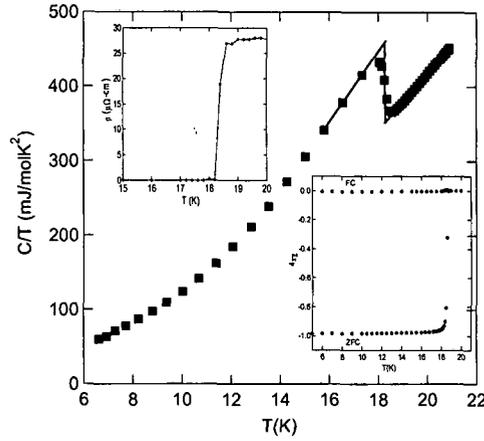
## 1. Introduction

Plutonium is a fascinating metal whose 5f electrons are poised on the boundary between localized and itinerant behavior. This instability gives rise to an extremely complex metallurgy[1] and challenges the state of the art in electronic structure calculations[2]. The crossover from localized to itinerant f-electron behavior is also central to the phenomenology of heavy fermion compounds[3].

Here, we discuss a recently discovered microcosm of the fascinating properties of plutonium: the discovery of superconductivity in  $\text{PuCoGa}_5$  at 18.5 K[4]. Not only is this a rather high  $T_c$  for an intermetallic compound but also there is at least the suggestion that this superconductivity may be unconventional and, perhaps, spin-fluctuation mediated[5, 6].

## 2. Evidence for superconductivity in $\text{PuCoGa}_5$

Large single crystals of  $\text{PuCoGa}_5$  have been grown from an excess Ga flux. Further and independently, single-crystal platelets have been obtained by arc-melting and subsequent annealing. Physical properties of these materials are identical and reveal bulk superconductivity near 18.5 K[4]. In both cases, single-crystal structural determinations have been made. One finds that  $\text{PuCoGa}_5$  crystallizes in the  $\text{HoCoGa}_5$  crystal structure with tetragonal lattice constants  $a=4.232$  Å and  $c=6.786$  Å. This is the same crystal structure in which  $\text{CeMIn}_5$  ( $M=\text{Co, Rh, Ir}$ ), a family of unconventional heavy fermion superconductors[7], and  $\text{UMGa}_5$ [8] also crystallize.

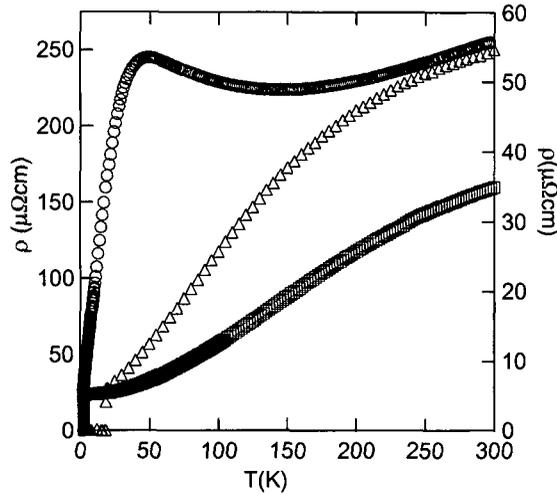


**Figure 1.** Evidence for superconductivity in  $\text{PuCoGa}_5$ . The main body of the figure shows heat capacity plotted as  $C/T$  versus  $T$ . The upper and lower insets show electrical resistivity and magnetic susceptibility, respectively, as a function of temperature. In all cases signatures of a phase transition are observed in the vicinity of 18.5 K.

In Figure 1 we show the evidence for bulk superconductivity in  $\text{PuCoGa}_5$ . A transition to zero resistance, coincident with full-shielding diamagnetism, is observed near 18.5 K. At this same temperature, a step-like transition in heat capacity is observed. If one assumes the BCS value of  $\Delta C/\gamma T_c = 1.43$ , then one infers from these data that  $\gamma$ , a measure of the conduction electron contribution to the low-temperature heat capacity, is  $77 \text{ mJ/molK}^2$ . This value of  $\gamma$  is enhanced relative to that expected for normal metals and is suggestive of heavy fermion behavior.

Interestingly, the  $T_c$  of  $\text{PuCoGa}_5$  decreases from its initial value of  $\sim 18.5$  K as a function of time at a rate of approximately 0.2 K/month. This decrease would seem to be a result of radiation-induced self-damage associated with the spontaneous decay of  $^{239}\text{Pu}$ . This mechanism is further indicated by the fact that the critical current,  $J_c$ , actually increases with time over the same period[4].

A correspondingly large value of the upper critical field  $H_{c2}$  in  $\text{PuCoGa}_5$  has been inferred[4]. In particular, field-dependent resistivity data yield an initial slope of  $dH_{c2}/dT$  of  $-59 \text{ kOe/K}$ . From this value, one can estimate an upper critical field of 740 kOe. Further, one can estimate the BCS coherence length and therefore the Fermi velocity, and find that  $\gamma \sim 60 \text{ mJ/molK}^2$  in the free-electron limit. Similarly,



**Figure 2.** Resistivity as a function of temperature for CeCoIn<sub>5</sub> (circles), PuCoGa<sub>5</sub> (triangles), and UCoGa<sub>5</sub> (squares). The data for CeCoIn<sub>5</sub> are plotted on the right axis, whereas PuCoGa<sub>5</sub> and UCoGa<sub>5</sub> use the left.

from estimates of the thermodynamic critical field, one can estimate  $\gamma \sim 70$  mJ/molK<sup>2</sup>, assuming the BCS value for the condensation energy. Thus, one has several independent estimates of  $\gamma \sim 100$  mJ/molK<sup>2</sup> in PuCoGa<sub>5</sub>. Although this is a rather small value compared to other heavy fermion superconductors, it is significantly enhanced compared to UMGa<sub>5</sub>[8], in which no superconductivity is observed.

### 3. Trends in normal state properties

Normal-state properties provide further evidence that PuCoGa<sub>5</sub> displays stronger electron correlation effects than UCoGa<sub>5</sub>. Figure 2 displays electrical resistivity for CeCoIn<sub>5</sub>, PuCoGa<sub>5</sub>, and UCoGa<sub>5</sub>. Judging from the characteristic change in curvature of the temperature dependence of the resistivity, one can see in Figure 2 that this temperature scale is higher in PuCoGa<sub>5</sub> than in CeCoIn<sub>5</sub>, but not as high as in UCoGa<sub>5</sub>. This trend can be confirmed from heat capacity measurements in these compounds, which find that the linear-in-temperature coefficient of the low-temperature heat capacity,  $\gamma$ , increases from  $\sim 10$  mJ/molK<sup>2</sup> in UCoGa<sub>5</sub> to  $\sim 100$  mJ/molK<sup>2</sup> in PuCoGa<sub>5</sub> to  $\sim 1000$  mJ/molK<sup>2</sup> in CeCoIn<sub>5</sub>[7].

From these data one is led to the conclusion that the superconductivity in PuCoGa<sub>5</sub> may be unconventional. In such a scenario, the order-of-magnitude higher  $T_c$  in PuCoGa<sub>5</sub> compared to CeCoIn<sub>5</sub> ( $T_c=2.3$  K)[7] would be expected from the increase in bandwidth in going from 5f electrons to 4f electrons[9, 10]. It is generally understood that 4f electrons have a greater degree of localization than do 5f electrons, as deduced, for example, by the evolution of the Wigner-Seitz radius as a function of atom across the lanthanide/actinide families[11].

Although the suggestion of unconventional superconductivity in PuCoGa<sub>5</sub> may seem implausible, the alternative, an 18-K conventional, phonon-mediated superconductor, is equally challenging. In PuCoGa<sub>5</sub>, magnetic susceptibility measurements reveal Curie-Weiss behavior consistent with a paramagnetic moment of  $\sim 0.7 \mu_B/\text{Pu}$ [4]. The pair-breaking tendency of magnetic moments would suggest that UCoGa<sub>5</sub>, a temperature-independent paramagnet, would have a higher  $T_c$  than PuCoGa<sub>5</sub>, in contrast to what is observed.

#### 4. Summary

We have discussed the observation of superconductivity above 18 K in PuCoGa<sub>5</sub>. The speculation that this might be unconventional, spin-fluctuation-mediated superconductivity raises the possibility that PuCoGa<sub>5</sub> could be an intellectual bridge between the known heavy-fermion superconductors (with characteristic  $T_c \sim 1$  K) and the high- $T_c$  cuprates (with characteristic  $T_c \sim 100$  K). Thus, the transuranics may represent a particularly fertile, if unplowed, field for the discovery of additional superconductors.

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